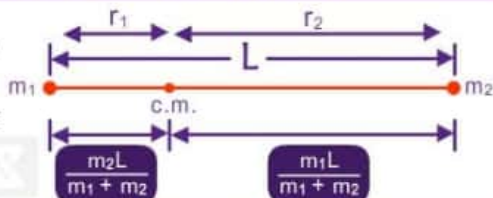


# CENTRE OF MASS OF SOME COMMON SYSTEM

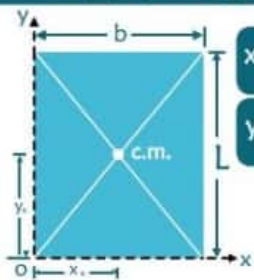
## System of Two Point Masses

$$m_1 r_1 = m_2 r_2$$

The Centre of mass lies closer to the heavier Mass.



## Rectangular Plate (By symmetry)

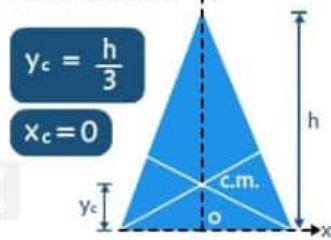


$$x_c = \frac{b}{2}$$

$$y_c = \frac{L}{2}$$

## Triangular Plate (By qualitative argument)

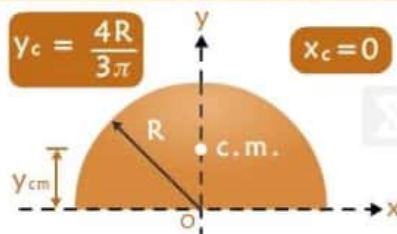
At the Centroid:



$$y_c = \frac{h}{3}$$

$$x_c = 0$$

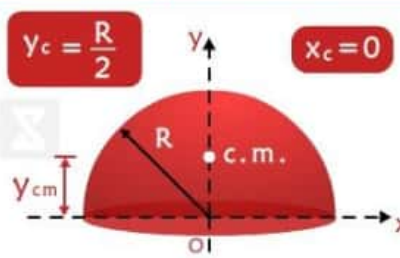
## Semi-Circular Disc



$$y_c = \frac{4R}{3\pi}$$

$$x_c = 0$$

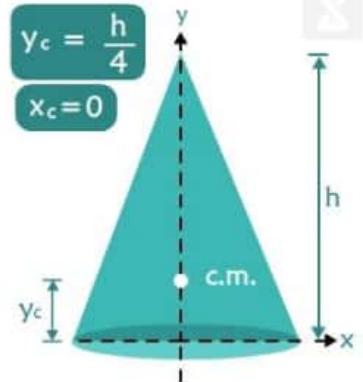
## Hemispherical Shell



$$y_c = \frac{R}{2}$$

$$x_c = 0$$

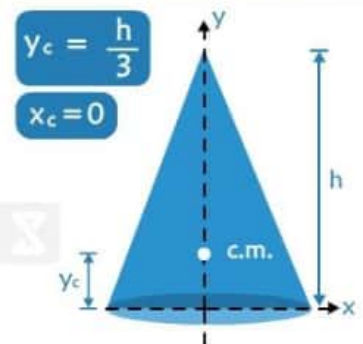
## Circular Cone (Solid)



$$y_c = \frac{h}{4}$$

$$x_c = 0$$

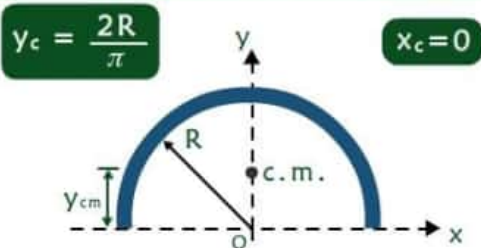
## Circular Cone (Hollow)



$$y_c = \frac{h}{3}$$

$$x_c = 0$$

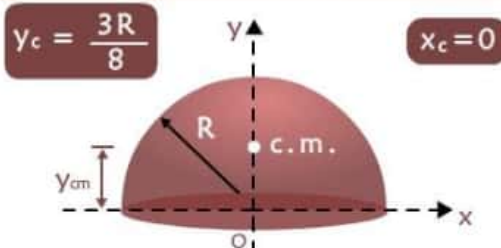
## Semi-Circular Ring



$$y_c = \frac{2R}{\pi}$$

$$x_c = 0$$

## Solid Hemisphere



$$y_c = \frac{3R}{8}$$

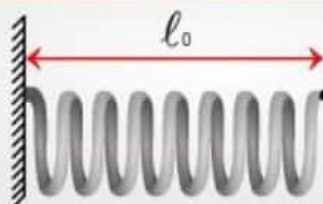
$$x_c = 0$$

# SPRING FORCE

1

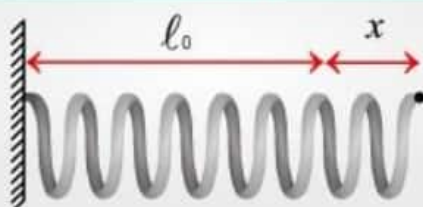
## STRETCHED SPRING

Initial length ( $\ell_0$ )



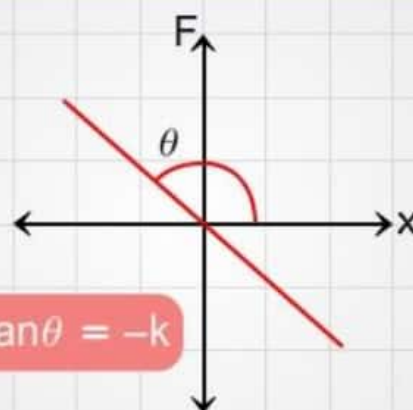
$$F = 0$$

Stretched by  $x$



$$F = -kx$$

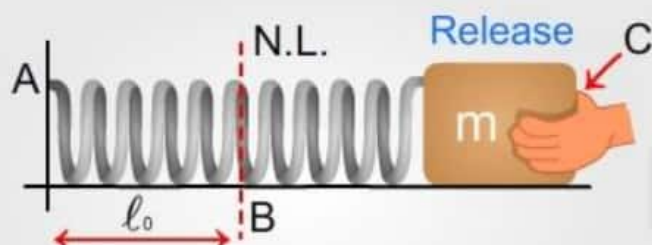
Spring Force v/s Displacement



2

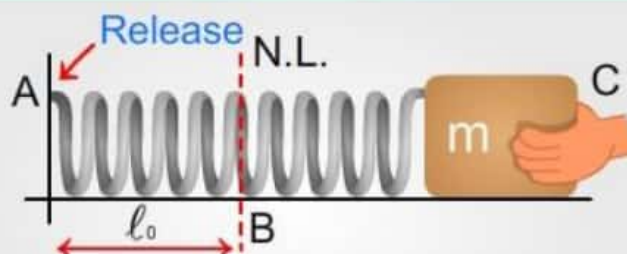
## SPRING ATTACHED TO A BLOCK

Released at C



When the block is released at point C then spring force doesn't change instantaneously because of friction at mass m.

Released at A



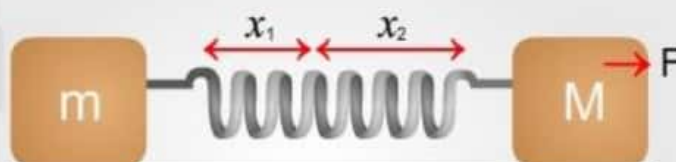
When point A is released then the spring force changes instantaneously to become zero.

3

## SPRING BLOCK SYSTEM



Maximum Extension  $x_{\max} = v_0 \sqrt{\frac{2}{3k} m}$



$$x_{\max} = x_1 + x_2 = \frac{2mF}{k(m + M)}$$



# IMPULSE AND MOMENTUM

## IMPULSE

Impulse of a force ' $F$ ' acting on a body for a time interval  $t = t_1$  to  $t = t_2$  is defined as

$$\vec{I} = \int_{t_1}^{t_2} \vec{F} dt$$

$$\vec{I}_{Re} = \int_{t_1}^{t_2} \vec{F}_{Res} dt = \Delta \vec{P}$$

(Impulse - Momentum Theorem)

## COEFFICIENT OF RESTITUTION ( $e$ )

The coefficient of restitution is defined as the ratio of the impulses of reformation and deformation of either body.

$$e = \frac{\text{Impulse of reformation}}{\text{Impulse of deformation}} = \frac{\int F_r dt}{\int F_d dt}$$

$$e = \frac{\text{Velocity of separation of point of contact}}{\text{Velocity of approach of point of contact}}$$

## LINEAR MOMENTUM

Linear momentum is a vector quantity defined as the product of an object's mass  $m$ , and its velocity  $v$ . Linear momentum is denoted by the letter  $p$  and is called "momentum" in short:

$$p = mv$$

Note that a body's momentum is always in the same direction as its velocity vector. The units of momentum are  $\text{kg.m/s}$ .

## CONSERVATION OF LINEAR MOMENTUM

For a single mass or single body, If net force acting on the body is zero. Then,

$$\vec{p} = \text{constant} \quad \text{or} \quad \vec{v} = \text{constant}$$

(if mass = constant)

If net external force acting on a system of particles or system of rigid bodies is zero. Then,

$$\vec{P}_{CM} = \text{constant} \quad \text{or} \quad \vec{V}_{CM} = \text{constant}$$



# COLLISION



**Note :-** In every type of collision, only linear momentum remains constant.

## HEAD ON ELASTIC COLLISION



Before Collision



After Collision

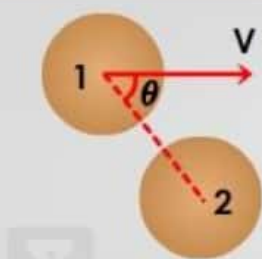
In this case, linear momentum and kinetic energy both are conserved. After solving two conservation equations. We get,

$$v'_1 = \left( \frac{m_1 - m_2}{m_1 + m_2} \right) v_1 + \left( \frac{2m_2}{m_1 + m_2} \right) v_2 \quad \text{and} \quad v'_2 = \left( \frac{m_2 - m_1}{m_1 + m_2} \right) v_2 + \left( \frac{2m_1}{m_1 + m_2} \right) v_1$$

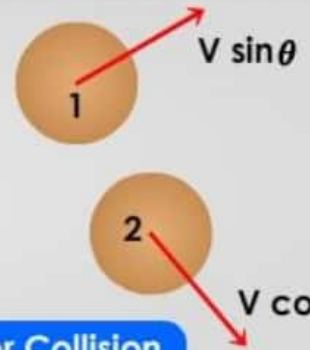
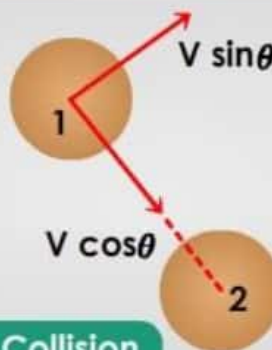
## HEAD ON INELASTIC COLLISION

- ➔ In an inelastic collision, the colliding particles do not regain their shape and size completely after the collision.
- ➔ Some fraction of mechanical energy is retained by the colliding particles in the form of deformation potential energy. Thus, the kinetic energy of the particles no longer remains conserved.
- ➔  $(\text{Energy loss})_{\text{Perfectly Inelastic}} > (\text{Energy loss})_{\text{Partial Inelastic}}$
- ➔  $0 < e < 1$  :  $e$  = coefficient of restitution

## OBLIQUE COLLISION (BOTH ELASTIC IN ELASTIC)



Before Collision



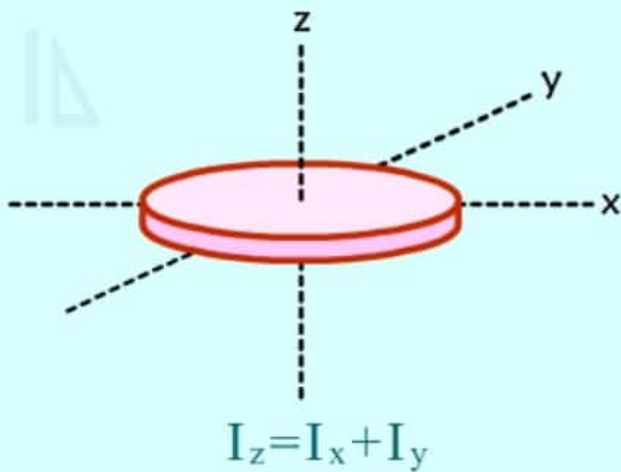
After Collision

BALL	COMPONENT ALONG COMMON TANGENT DIRECTION		COMPONENT ALONG COMMON NORMAL DIRECTION	
	Before Collision	After Collision	Before Collision	After Collision
1	$V \sin \theta$	$V \sin \theta$	$V \cos \theta$	0
2	0	0	0	$V \cos \theta$

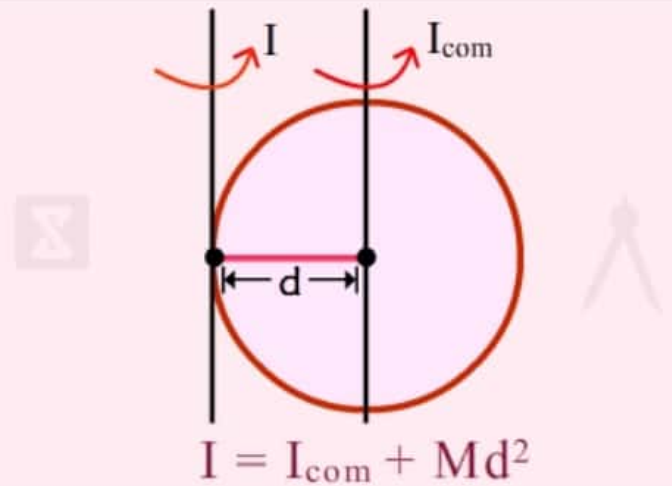




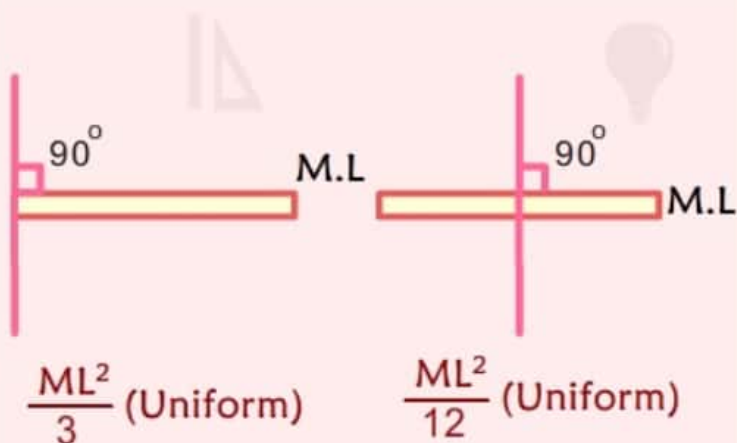
## Perpendicular Axis Theorem



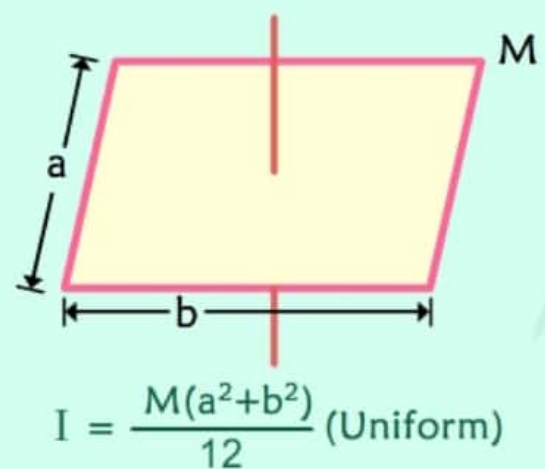
## Parallel Axis Theorem



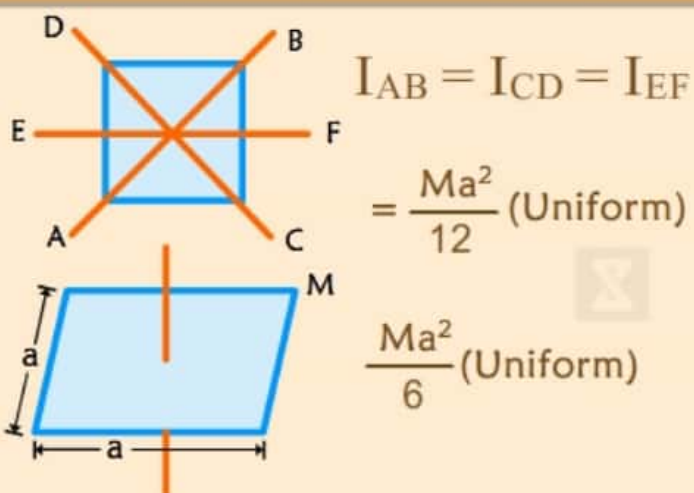
## Rod



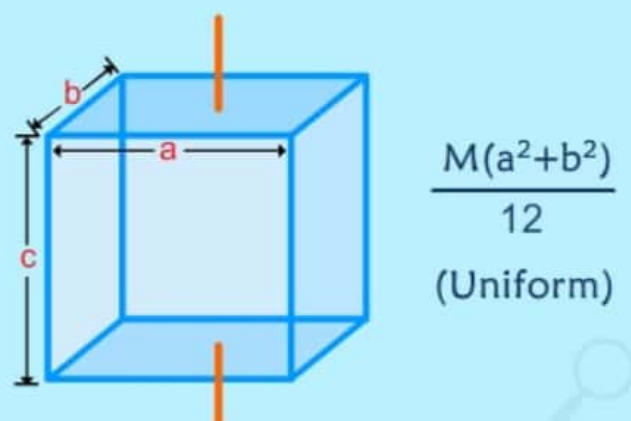
## Rectangular Plate



## Square Plate



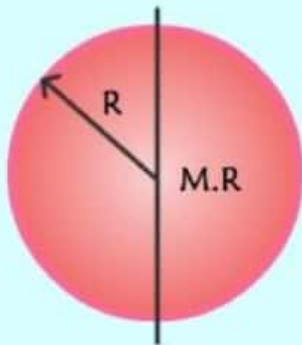
## Cuboid



# MOMENT OF INERTIA

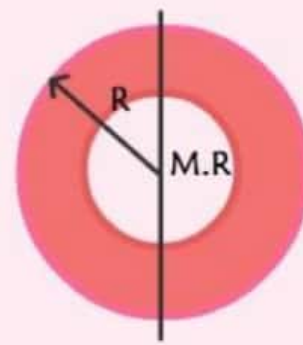
Part II

## Solid Sphere



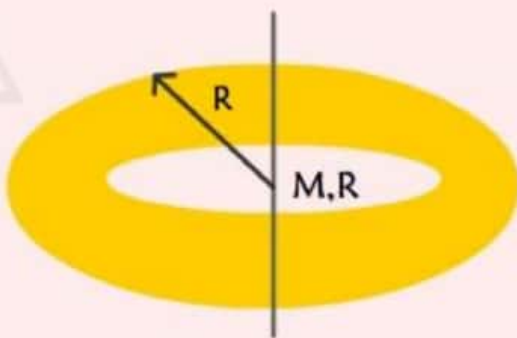
$$I = \frac{2}{5} MR^2 \text{ (Uniform)}$$

## Hollow Sphere



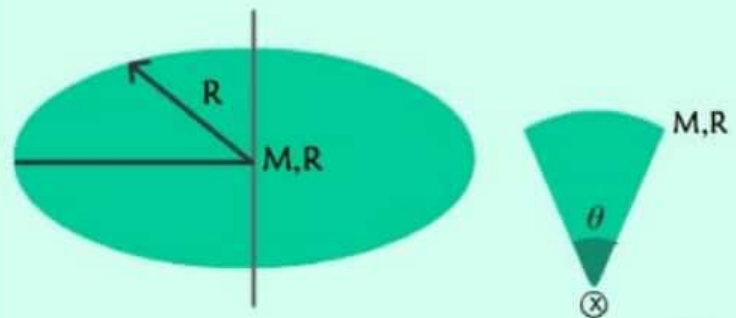
$$I = \frac{2}{3} MR^2 \text{ (Uniform)}$$

## Ring



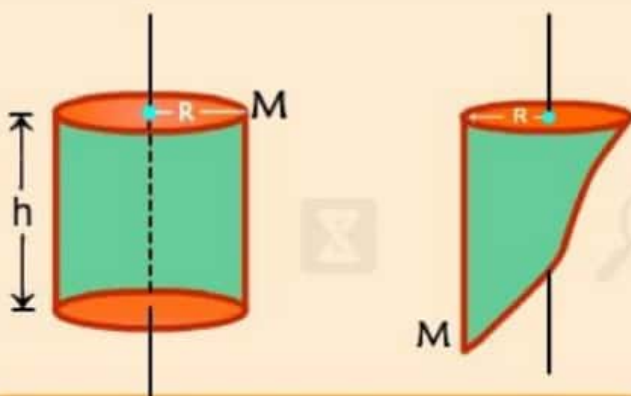
$$I = MR^2 \text{ (Uniform or Non Uniform)}$$

## Disc



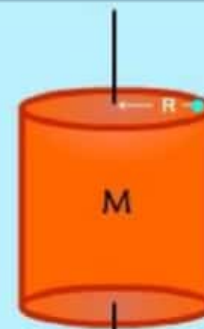
$$I = \frac{MR^2}{2} \text{ (Uniform)}$$

## Hollow cylinder



$$I = MR^2 \text{ (Uniform or Non Uniform)}$$

## Solid cylinder



$$I = \frac{MR^2}{2} \text{ (Uniform)}$$

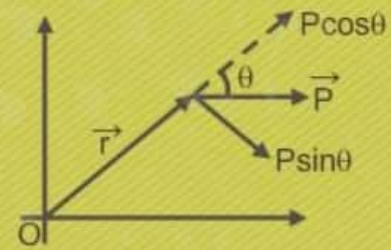


# ANGULAR MOMENTUM



## 1 ANGULAR MOMENTUM OF A PARTICLE ABOUT A POINT

$$\vec{L} = \vec{r} \times \vec{P} \Rightarrow L = rP \sin\theta$$



## 2 ANGULAR MOMENTUM OF A RIGID BODY ROTATING ABOUT A FIXED AXIS

$$L = I\omega$$

Here,  $I$  is the moment of inertia of the rigid body about axis.

## 3 CONSERVATION OF ANGULAR MOMENTUM

The law of conservation of angular momentum states that when **no external torque acts** on an object, **no change of angular momentum** will occur.

Since  $\vec{\tau}_{\text{net}} = \frac{d\vec{L}}{dt}$ . Now if,  $\vec{\tau}_{\text{net}} = 0$ , then  $\frac{d\vec{L}}{dt} = 0$ , so that  $\vec{L} = \text{constant}$ .

## 4 ANGULAR IMPULSE

The angular impulse of a torque in a given time interval is defined as

$$\int_{t_1}^{t_2} \vec{\tau} dt$$

$$\vec{\tau} = \frac{d\vec{L}}{dt}$$

## UNIFORM PURE ROLLING

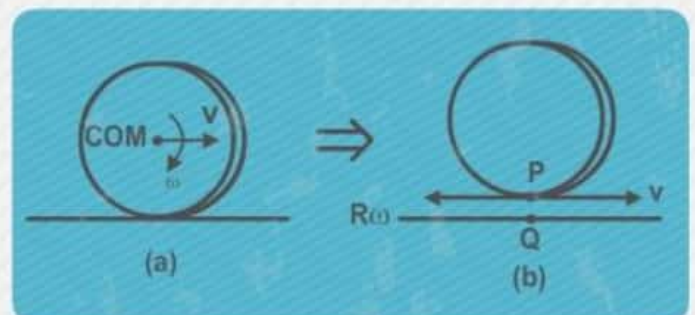
Pure rolling means no relative motion (or no slipping at point of contact between two bodies.)

$$V_P = V_Q \quad \text{or} \quad V - R\omega = 0 \quad \text{or} \quad V = R\omega$$

If  $V_P > V_Q$  or  $V > R\omega$ , the motion is said to be forward slipping and if

$V_P < V_Q < R\omega$ , the motion is said to be backward slipping.

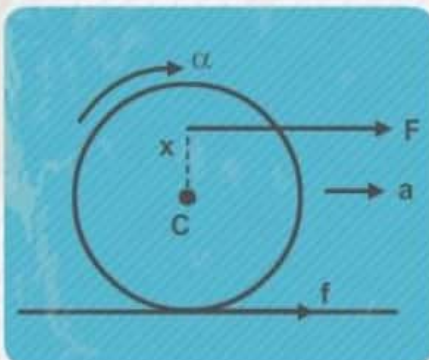
The condition of pure rolling on a stationary ground is,  
 $a = R\alpha$





## 1 PURE ROLLING WHEN FORCE F ACT ON A BODY

Suppose a force  $F$  is applied at a distance  $x$  above the centre of a rigid body of radius  $R$ , mass  $M$  and moment of inertia  $CMR^2$  about an axis passing through the centre of mass. Applied force  $F$  can produces by itself a linear acceleration  $a$  and an angular acceleration  $\alpha$ .



$a$  = liner acceleration ,  $\alpha$  = angular acceleration from linear motion

$$F + f = Ma$$

From rotational motion :  $Fx - fR = I\alpha$

$$a = \frac{F(R+x)}{MR(C+1)}, \quad f = \frac{F(x-RC)}{R(C+1)}$$

## 2 PURE ROLLING ON A INCLINED PLANS

A rigid body of radius  $R$ , and mass  $m$  is released at rest from height  $h$  on the incline whose inclination with horizontal is  $\theta$  and assume that friciton is sufficient for pure rolling then,

$$a = \alpha R \text{ and } v = R\omega$$

$\omega$  = Angular Velocity

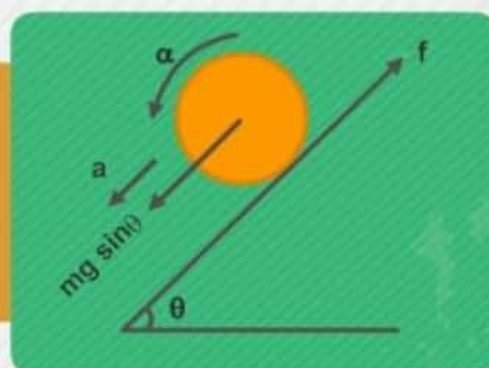
$\alpha$  = Angular Acceleration

Linear Acceleration,

$$\alpha = \frac{g \sin \theta}{1 + C}$$

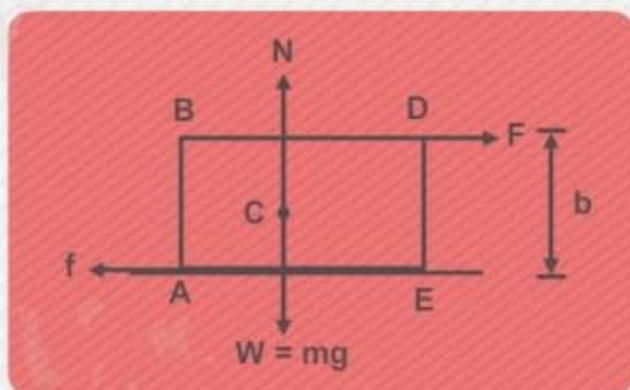
$C$  = Center of Mass

So, body which have low value of  $C$  have greater acceleration.



## TOPPLING

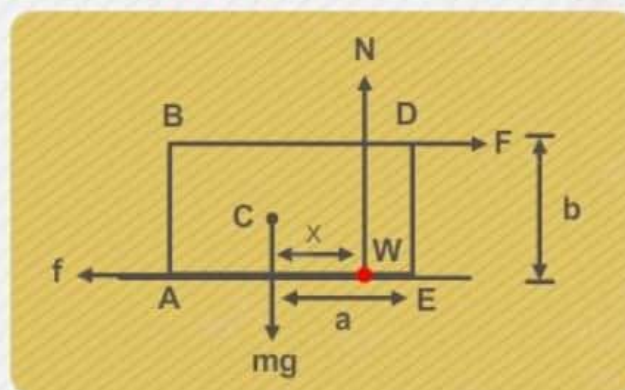
### Torque about E



Balancing Torque at E

$$Fb = (mg)a \implies a = \frac{Fb}{mg}$$

### Torque about W



Balancing Torque at W

$$Fb + N(a-x) = mga$$

if  $x = a$

$$F_{\max} b = mga \implies F_{\max} = \frac{mga}{b}$$

